

electric vector at an angle  $\varphi$  to the reflecting plane. After reflexion, a ray proceeds along  $MO$  with components of amplitude proportional to  $\sin \varphi \cos 2\alpha$  and  $\cos \varphi$ , and polarized parallel to  $OX$  and  $OZ$  respectively. After diffraction at  $O$ , the  $X$  component will give rise to components of amplitude proportional to

$-\sin \varphi \cos 2\alpha \sin \chi \sin \gamma_0$  polarized in the plane  $POZ$  and  $\sin \varphi \cos 2\alpha \cos \gamma_0$  polarized perpendicular to this plane. The  $Z$  component gives corresponding components of amplitude proportional to  $\cos \varphi \cos \chi$  polarized in the plane  $POZ$  and zero in the perpendicular direction. The intensity diffracted along  $OP$  is accordingly proportional to

$$(\cos \varphi \cos \chi - \sin \varphi \sin \chi \sin \gamma_0 \cos 2\alpha)^2 + \sin^2 \varphi \cos^2 2\alpha \cos^2 \gamma_0.$$

After averaging this result over all values of  $\varphi$ , to allow for the unpolarized nature of the rays incident at  $M$ , we obtain

$$\frac{1}{2}(\cos^2 \chi + \sin^2 \chi \sin^2 \gamma_0 \cos^2 2\alpha + \cos^2 \gamma_0 \cos^2 2\alpha).$$

Since, by a well known result, the intensity incident at  $O$  is proportional to  $\frac{1}{2}(1 + \cos^2 2\alpha)$ , the polarization factor is

$$P = \frac{\cos^2 \chi + \sin^2 \chi \sin^2 \gamma_0 \cos^2 2\alpha + \cos^2 \gamma_0 \cos^2 2\alpha}{1 + \cos^2 2\alpha}. \quad (3)$$

The angles  $\gamma_0$  and  $\chi$  are suitable film co-ordinates only in the normal-beam arrangement. In the inclined-beam arrangement suitable co-ordinates are  $\mu$ ,  $\nu$  and  $\gamma$ , where  $\mu$ ,  $\nu$  are the inclinations of the incident and diffracted beams to the equatorial plane of the camera, and  $\gamma$  is the azimuth of the diffracted ray defined in the usual way (Buerger, 1942, p. 297). It can then be shown by appropriate transformation of the axes of co-ordinates that

$$\cos \chi \sin \gamma_0 = \cos \nu \sin \gamma$$

and

$$\cos^2 \chi = \cos^2 \nu \sin^2 \gamma + \cos^2 \mu \cos^2 \nu \cos^2 \gamma + \sin^2 \mu \sin^2 \nu + \frac{1}{2} \sin 2\mu \sin 2\nu \cos \gamma;$$

and hence, by substitution and rearrangement,

$$P = \{\cos^2 2\alpha + \cos^2 \mu \cos^2 \nu + \sin^2 \mu \sin^2 \nu + \cos^2 \nu \sin^2 \gamma (\sin^2 2\alpha - \cos^2 \mu) + \frac{1}{2} \sin 2\mu \sin 2\nu \cos \gamma\} / \{1 + \cos^2 2\alpha\}.$$

For purposes of numerical calculation this expression can be regarded as being of the form

$$P = A + B \sin^2 \gamma + C \cos \gamma, \quad (5)$$

where  $A$ ,  $B$  and  $C$  are constants for any given layer line recorded with a given camera setting.

By putting  $\alpha = 0$  in the above expression we obtain

$$P = \frac{1}{2}\{1 + (\cos \mu \cos \nu \cos \gamma + \sin \mu \sin \nu)^2\} \quad (6)$$

as a convenient form of the factor for unpolarized radiation in terms of the same co-ordinates. Expression (6) can of course be more simply derived by direct trigonometrical transformation from (1).

I wish to thank the Directors of Ferodo Ltd for permission to publish this communication.

### References

- BUERGER, M. J. (1942). *X-Ray Crystallography*. New York: Wiley.  
*Internationale Tabellen zur Bestimmung von Kristallstrukturen* (1935). Berlin: Borntraeger.

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